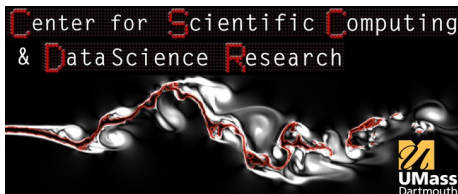


Strong Stability Preserving Two Derivative Methods

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In the numerical solution of hyperbolic partial differential equations

$$u_t + f(u)_x = 0$$

with discontinuous solutions, we need the spatial discretization to satisfy certain nonlinear stability properties usually **when coupled with forward Euler**.

Strong stability preserving (SSP) high order time discretizations were developed to ensure that these stability properties of the spatial discretization are preserved **when coupled with these higher order methods**.

The forward Euler condition

Given a system of ODEs, generally resulting from a spatial discretization of a PDE, of the form

$$u_t = \mathbf{F}(u)$$

Forward Euler condition:

$$\|u^n + \Delta t \mathbf{F}(u^n)\| \leq \|u^n\| \quad \forall \Delta t \leq \Delta t_{\text{FE}}$$

where $\|\cdot\|$ is some convex functional.

SSP methods extend this strong stability to higher order methods. How?

High order SSP Runge–Kutta methods

Consider the two-stage second order Runge–Kutta method

$$\begin{aligned}u^{(1)} &= u^n + \Delta t \mathbf{F}(u^n) \\ u^{n+1} &= \frac{1}{2}u^n + \frac{1}{2} \left(u^{(1)} + \Delta t \mathbf{F}(u^{(1)}) \right),\end{aligned}$$

we can write this method as a convex combination of the forward Euler method

$$\begin{aligned}\|u^{(1)}\| &= \|u^n + \Delta t \mathbf{F}(u^n)\| \leq \|u^n\| \text{ for } \Delta t \leq \Delta t_{\text{FE}} \\ \|u^{n+1}\| &= \left\| \frac{1}{2}u^n + \frac{1}{2} \left(u^{(1)} + \Delta t \mathbf{F}(u^{(1)}) \right) \right\| \\ &\leq \frac{1}{2} \|u^n\| + \frac{1}{2} \|u^{(1)} + \Delta t \mathbf{F}(u^{(1)})\| \leq \|u^n\| \text{ for } \Delta t \leq \Delta t_{\text{FE}}\end{aligned}$$

This method is SSP with $\mathcal{C} = 1$ and $\mathcal{C}_{\text{eff}} = \frac{1}{2}$.

High order SSP Runge–Kutta methods

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This method is SSP with $\mathcal{C} = 1$ and $\mathcal{C}_{\text{eff}} = \frac{1}{2}$.

Can I always do this? No!

IF

we can decompose a higher order method into **convex combinations of forward Euler**

THEN

the Runge–Kutta or multistep method will preserve the strong stability property that F satisfies when coupled with FE.

SSP methods were introduced by Chi-Wang Shu and Stanley Osher in 1988 for time-evolution of hyperbolic PDEs with shocks. These were initially called TVD time stepping methods.

SSP Runge–Kutta Methods

- Explicit SSP Runge–Kutta methods have order $p \leq 4$ and $\mathcal{C}_{\text{eff}} \leq 1$.
- Explicit SSP Runge–Kutta methods for linear problems have no order barrier but $\mathcal{C}_{\text{eff}} \leq 1$.
- Implicit SSP Runge–Kutta methods have order $p \leq 6$
- Implicit Runge–Kutta cannot be **unconditionally** SSP.
- All optimal implicit SSP Runge–Kutta methods found have $\mathcal{C}_{\text{eff}} \leq 2$.

We try to find ways to break the order barriers and in particular the SSP coefficient bounds.

Idea: The key to the SSP property was to decompose the higher order methods into a convex combination of forward Euler steps, i.e.

$$\alpha \mathbf{u}^{(j)} + \Delta t \beta \mathbf{F}(\mathbf{u}^{(j)})$$

where α and β are non-negative.

But what if some of the β coefficients are negative?

Downwinding

If we have a spatial discretization $\tilde{\mathbf{F}}$ that approximates the same derivative as the spatial approximation \mathbf{F} , but has the property that

$$\|\mathbf{u}^n - \Delta t \tilde{\mathbf{F}}(\mathbf{u}^n)\| \leq \|\mathbf{u}^n\| \quad \text{for } \Delta t \leq \Delta t_{\text{FE}},$$

while \mathbf{F} has the property

$$\|\mathbf{u}^n + \Delta t \mathbf{F}(\mathbf{u}^n)\| \leq \|\mathbf{u}^n\| \quad \text{for } \Delta t \leq \Delta t_{\text{FE}}.$$

If a method can be decomposed into convex combination of $\mathbf{u}^n + \Delta t \mathbf{F}(\mathbf{u}^n)$ and $\mathbf{u}^n - \Delta t \tilde{\mathbf{F}}(\mathbf{u}^n)$ then the Runge–Kutta method will preserve the strong stability property that \mathbf{F} and $\tilde{\mathbf{F}}$ satisfy when coupled with forward Euler or backward Euler (respectively).

Downwinding: Breaking the order barrier

Explicit Runge–Kutta methods with downwinding can **break the order barrier**

If you have a PDE that admits an upwind stable discretization F and a downwind stable discretization \tilde{F} such that we can obtain

- A 4-stage fourth order method with $\mathcal{C} = 0.936$ (Gottlieb & Shu 1998)
- A fifth order explicit Runge–Kutta method with $\mathcal{C} = 1.17$ (Ruuth & Spiteri 2004)
- This may be more efficient than it first appear as we can compute the downwind and upwind operators together at less than double the cost (Gottlieb & Ruuth 2006).

Ketcheson (2011): Implicit Runge–Kutta methods with downwinding can **break the SSP coefficient bound**

$$y_1 = \frac{2}{r(r-2)}u^n + \frac{2}{r} \left(y_1 + \frac{\Delta t}{r} F(y_1) \right) + \frac{r^2 - 4r + 2}{r(r-2)} \left(y_2 - \frac{\Delta t}{r} \tilde{F}(y_2) \right)$$

$$y_2 = y_1 + \frac{\Delta t}{r} F(y_1)$$

$$u^{n+1} = y_2 + \frac{\Delta t}{r} F(y_2)$$

Similarly, a family of third order fully implicit downwind methods with arbitrarily large \mathcal{C} was found by Ketcheson's student Yiannis Hadjimichael and appears in his thesis (2017).

Another approach: using additional derivatives

To break the order barriers and in particular the SSP coefficient bounds we can also use an additional derivative.

This approach is based on a Taylor series idea, where $u_{tt} = \dot{\mathbf{F}}$ is used. In our case we add this to the Runge–Kutta method.

- Enhancing Runge–Kutta methods with additional derivatives as far back as 1950 (Turan) and 1963 (Stroud & Stroud).
- Multistage multiderivative time integrators for ordinary differential equations were studied by Shintani (1971,1972) and Kastlunger and G. Wanner (1972) and Mitsui (1982), Ono (2004) and Tsai (2010).
- More recently, for the the time evolution of partial differential equations, these were studied by Seal & Christlieb, Tsai, Li & Du. This makes sense for hyperbolic conservation laws, where it is natural (i.e. cheap) to include the second derivative.

Two derivatives Runge–Kutta methods

In 2014 Seal and Christlieb studied high-order multiderivative time integrators for hyperbolic conservation laws. In particular they used the 4th order method:

$$y^{(1)} = u^n + \frac{\Delta t}{2} \mathbf{F}(u^n) + \frac{\Delta t^2}{8} \dot{\mathbf{F}}(u^n)$$
$$u^{n+1} = u^n + \Delta t \mathbf{F}(u^n) + \frac{\Delta t^2}{6} (\dot{\mathbf{F}}(u^n) + 2\dot{\mathbf{F}}(y^{(1)})).$$

Is this method SSP?

Two derivative Runge–Kutta methods

We proved that the fourth order method above gives

$$\|u^{n+1}\| \leq \|u^n\|$$

if \mathbf{F} satisfies the forward Euler condition

Forward Euler

$$\|\mathbf{u}^n + \Delta t \mathbf{F}(\mathbf{u}^n)\| \leq \|\mathbf{u}^n\| \quad \text{for } \Delta t \leq \Delta t_{\text{FE}}.$$

and $\dot{\mathbf{F}}$ satisfies the condition

Second derivative

$$\|u^n + \Delta t^2 \dot{\mathbf{F}}(u^n)\| \leq \|u^n\| \quad \text{for } \Delta t \leq \Delta t_{\text{SD}} = K \Delta t_{\text{FE}}$$

where $K > 0$.

These building blocks make sense

Given the simple linear one-way wave equation $U_t = U_x$, let

$$u_t \approx u_x \approx \mathbf{F}(u^n)_j := \frac{1}{\Delta x} (u_{j+1}^n - u_j^n).$$

and

$$u_{tt} = u_{xx} \approx \dot{\mathbf{F}}(u^n)_j := \frac{1}{\Delta x^2} (u_{j+1}^n - 2u_j^n + u_{j-1}^n).$$

Then F and \dot{F} satisfy the total variation diminishing (TVD) property:

$$\|u^n + \Delta t \mathbf{F}(u^n)\|_{TV} \leq \|u^n\|_{TV} \quad \text{for } \Delta t \leq \Delta x,$$

$$\|u^n + \Delta t^2 \dot{\mathbf{F}}(u^n)\|_{TV} \leq \|u^n\|_{TV} \quad \text{for } \Delta t \leq \frac{\sqrt{2}}{2} \Delta x.$$

Optimal one stage second order methods

The Taylor series method is the unique explicit one-stage two-derivative second order method:

$$u^{n+1} = u^n + \Delta t F(u^n) + \frac{1}{2} \Delta t^2 \dot{F}(u^n).$$

It is one method, but the optimal Shu-Osher decomposition and the corresponding SSP coefficient, depend on the value of K in the second derivative condition.

Optimal one stage second order methods:

This method can be written as a convex combination of the building blocks:

$$u^{n+1} = (1 - \alpha) \left(u^n + \frac{1}{1 - \alpha} \Delta t F(u^n) \right) + \alpha \left(u^n + \frac{1}{2\alpha} \Delta t^2 \dot{F}(u^n) \right).$$

This is SSP for $\Delta t \leq \max\{(1 - \alpha)\Delta t_{FE}, \sqrt{2\alpha}K\Delta t_{FE}\}$, so we set

$$(1 - \alpha)^2 = 2\alpha K^2 \quad \implies \quad \alpha = 1 + K^2 \pm K\sqrt{2 + K^2}$$

to obtain

$$C = K\sqrt{2 + K^2} - K^2.$$

Two-stage third order methods

Optimal SSP explicit two-stage two-derivative third order methods take the form

$$\begin{aligned}u^* &= u^n + a\Delta t F(u^n) + \hat{a}\Delta t^2 \dot{F}(u^n), \\u^{n+1} &= u^n + b_1\Delta t F(u^n) + b_2\Delta t F(u^*) \\&\quad + \hat{b}_1\Delta t^2 \dot{F}(u^n) + \hat{b}_2\Delta t^2 \dot{F}(u^*).\end{aligned}$$

The optimal method has SSP coefficient $\mathcal{C} = r$ that depends on K . This is a little messy!

Two-stage third order methods

The optimal method has SSP coefficient $\mathcal{C} = r$ given by the smallest positive root of the polynomial

$$2K(a_0 - 2K) + 4K^3a_0 - a_0r + \frac{1 - a_0}{2K^2}r^2 - \frac{\frac{a_0}{2K} + K}{6K^3}r^3,$$

for $a_0 = \sqrt{K^2 + 2} - K$. The coefficients of the optimal method for each K depend on K and r and are given by

$$\begin{aligned} a &= \frac{1}{r} (Ka_0), & b_1 &= 1 - b_2, & b_2 &= \frac{2K^2(1 - \frac{1}{r}) + r}{Ka_0 + 2K^2} - \frac{r^2}{3K^2}, \\ \hat{a} &= \frac{1}{2}a^2, & \hat{b}_1 &= \frac{1}{2} - \frac{1}{2}ab_2 - \frac{1}{6a}, & \hat{b}_2 &= \frac{1}{6a} - \frac{1}{2}ab_2. \end{aligned}$$

Fourth order methods

The two-stage two-derivative fourth order method is

$$y^{(1)} = u^n + \frac{\Delta t}{2} \mathbf{F}(u^n) + \frac{\Delta t^2}{8} \dot{\mathbf{F}}(u^n)$$
$$u^{n+1} = u^n + \Delta t \mathbf{F}(u^n) + \frac{\Delta t^2}{6} (\dot{\mathbf{F}}(u^n) + 2\dot{\mathbf{F}}(y^{(1)})).$$

Although the method is unique, the optimal decomposition, and therefore the SSP coefficient, depends on K . The SSP coefficient $\mathcal{C} = r$ is given by the smallest positive root of

$$r^4 + 4K^2r^3 - 12K^2r^2 - 24K^4r + 24K^4.$$

Three-stage fifth order methods

Increasing the number of stages to three breaks the order barrier and allows fifth order method:

$$\begin{aligned}y^{(1)} &= u^n + a_{21}\Delta t F(u^n) + \dot{a}_{21}\Delta t^2 \dot{F}(u^n) \\y^{(2)} &= u^n + a_{31}\Delta t F(u^n) + \dot{a}_{31}\Delta t^2 \dot{F}(u^n) + \dot{a}_{32}\Delta t^2 \dot{F}(y^{(1)}) \\u^{n+1} &= u^n + \Delta t F(u^n) + \Delta t^2 \left(\dot{b}_1 \dot{F}(u^n) + \dot{b}_2 \dot{F}(y^{(1)}) + \dot{b}_3 \dot{F}(y^{(2)}) \right).\end{aligned}$$

The SSP coefficient $\mathcal{C} = r$ depends on K and is messy.

Three-stage fifth order methods

The SSP coefficient $\mathcal{C} = r$ is the largest positive root of

$$10r^2 a_{21}^4 - (100K^2 + 10r^2)a_{21}^3 + (130K^2 + 3r^2)a_{21}^2 - 50K^2 a_{21} + 6K^2,$$

where

$$a_{21} = \frac{K^6}{r^6} \left(-\frac{2}{K^4} r^5 + \frac{10}{K^4} r^4 + \frac{40}{K^2} r^3 - \frac{120}{K^2} r^2 - 240r + 240 \right).$$

Given K , we can find the corresponding r , and the coefficients are then given as a one-parameter system

$$\begin{aligned} \dot{a}_{21} &= \frac{1}{2} a_{21}^2, & a_{31} &= \frac{3/5 - a_{21}}{1 - 2a_{21}}, \\ \dot{a}_{32} &= \frac{1}{10} \left(\frac{(\frac{3}{5} - a_{21})^2}{a_{21}(1 - 2a_{21})^3} - \frac{\frac{3}{5} - a_{21}}{(1 - 2a_{21})^2} \right), & \dot{a}_{31} &= \frac{1}{2} \frac{(\frac{3}{5} - a_{21})^2}{(1 - 2a_{21})^2} - \dot{a}_{32}, \\ \dot{b}_2 &= \frac{2a_{31} - 1}{12a_{21}(a_{31} - a_{21})}, & \dot{b}_3 &= \frac{1 - 2a_{21}}{12a_{31}(a_{31} - a_{21})}, & \dot{b}_1 &= \frac{1}{2} - \dot{b}_2 - \dot{b}_3. \end{aligned}$$



Two derivative Runge–Kutta methods

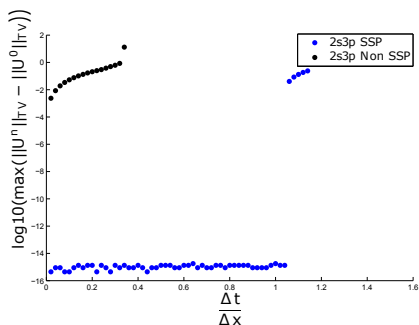
- Decomposing methods into convex combinations of the building blocks $u^n + \Delta t \mathbf{F}(u^n)$ and $u^n + \Delta t^2 \dot{\mathbf{F}}(u^n)$ allows us to build SSP two derivative Runge–Kutta methods.
- The SSP coefficient depends not only on the number of stages but also on K
- Using these building blocks we found explicit optimal SSP methods up to sixth order
- Two derivative methods allow us to break the order barrier of explicit Runge–Kutta methods
- Does the SSP property really matter in practice?

Yes!

Using an SSP and non-SSP two-stage two-derivative third order method on the motivating example above with a step function initial condition

$$u_0(x) = \begin{cases} 1 & \text{if } \frac{1}{4} \leq x \leq \frac{1}{2}, \\ 0 & \text{otherwise,} \end{cases}$$

and periodic boundary conditions $u(0, t) = u(1, t)$.



Maximal TV rise above the initial TV.

Different building blocks

Forward Euler

$$\|\mathbf{u}^n + \Delta t \mathbf{F}(\mathbf{u}^n)\| \leq \|\mathbf{u}^n\| \quad \text{for } \Delta t \leq \Delta t_{\text{FE}}.$$

&

Taylor Series condition

$$\|u^n + \Delta t \mathbf{F}(u^n) + \frac{1}{2} \Delta t^2 \dot{\mathbf{F}}(u^n)\| \leq \|u^n\| \quad \text{for } \Delta t \leq \Delta t_{\text{TS}}$$

↑ allows for **more** spatial discretizations

↓ fewer time-discretization methods are SSP (barrier of $p \leq 6$).

More flexibility in spatial discretization

For $U_t = U_x$ we define (as above)

$$\mathbf{F}(u^n)_j := \frac{1}{\Delta x} (u_{j+1}^n - u_j^n) \approx U_x(x_j).$$

To get $\dot{\mathbf{F}}$ apply this differentiation operator twice

$$\dot{\mathbf{F}}(u^n)_j := \frac{1}{\Delta x^2} (u_{j+2}^n - 2u_{j+1}^n + u_j^n).$$

\mathbf{F} satisfies the total variation diminishing (TVD) property for $\Delta t \leq \Delta x$, while the Taylor series term using \mathbf{F} and $\dot{\mathbf{F}}$ satisfies

$$\|u^n + \Delta t \mathbf{F}(u^n) + \frac{1}{2} \Delta t^2 \dot{\mathbf{F}}(u^n)\|_{TV} \leq \|u^n\|_{TV} \quad \text{for } \Delta t \leq \Delta x.$$

Less flexibility in time discretization

For example, the fourth order method

$$y^{(1)} = u^n + \frac{\Delta t}{2} F(u^n) + \frac{\Delta t^2}{8} \dot{F}(u^n)$$
$$u^{n+1} = u^n + \Delta t F(u^n) + \frac{\Delta t^2}{6} (\dot{F}(u^n) + 2\dot{F}(y^{(1)})).$$

cannot be written as convex combinations of these building blocks.
But this fourth order method **can**, with $C = 1$

$$y^{(1)} = u^n$$
$$y^{(2)} = u^n + \Delta t F(y^{(1)}) + \frac{1}{2} \Delta t^2 \dot{F}(y^{(1)})$$
$$y^{(3)} = u^n + \frac{1}{27} \Delta t (14F(y^{(1)}) + 4F(y^{(2)})) + \frac{2}{27} \Delta t^2 \dot{F}(y^{(1)})$$
$$u^{n+1} = u^n + \frac{1}{48} \Delta t (17F(y^{(1)}) + 4F(y^{(2)}) + 27F(y^{(3)})) + \frac{1}{24} \Delta t^2 \dot{F}(y^{(1)}).$$

Implicit methods with two derivatives

Whichever of these we use

Second derivative

$$u^n + \Delta t^2 \dot{F}(u^n)$$

Taylor Series

$$u^n + \Delta t F(u^n) + \frac{1}{2} \Delta t^2 \dot{F}(u^n)$$

when we move to implicit methods our time-step is very limited.
Why? Look at implicit Taylor series method

$$u^{n+1} = u^n + \Delta t F(u^{n+1}) - \frac{1}{2} \Delta t^2 \dot{F}(u^{n+1}).$$

Negative derivative condition

This suggests* the negative derivative condition

Negative derivative condition

$$\|u^n - \Delta t^2 \dot{F}(u^n)\| \leq \|u^n\| \quad \text{for } \Delta t \leq \Delta t_{\text{ND}}.$$

Or even

Implicit negative derivative condition

$$u^{n+1} = u^n - \Delta t^2 \dot{F}(u^{n+1}) \quad \implies \quad \|u^{n+1}\| \leq \|u^n\| \quad \forall \Delta t > 0.$$

↑ **Better SSP properties (unconditional)**

↓ **Fewer discretizations satisfy these conditions**

Unconditionally SSP implicit multiderivative schemes

For $u_t = \mathbf{G}(u)$, with $u_{tt} = \dot{\mathbf{G}}(u)$ we write the s -stage method in a special Shu-Osher form with only implicit computations

$$u^{(i)} = r_i u^n + \sum_{j=1}^{i-1} p_{ij} u^{(j)} + \Delta t d_{ii} G(u^{(i)}) + \Delta t^2 \dot{d}_{ii} \dot{G}(u^{(i)}),$$
$$u^{n+1} = u^{(s)}.$$

In matrix form, this becomes

$$U = \mathbf{R}e u^n + \mathbf{P}U + \Delta t \mathbf{D}G(U) + \Delta t^2 \dot{\mathbf{D}}\dot{G}(U),$$

where \mathbf{P} , and $\mathbf{R} = I - \mathbf{P}$ are $s \times s$ matrices, r_i are the i th row sum of \mathbf{R} , \mathbf{D} and $\dot{\mathbf{D}}$ are $s \times s$ diagonal matrices, and e is a vector of ones.

Theorem: Unconditionally SSP implicit schemes

Let the operators G and \dot{G} satisfy

Implicit Euler

$$u^{n+1} = u^n + \Delta t \mathbf{G}(u^{n+1}) \quad \implies \quad \|u^{n+1}\| \leq \|u^n\| \quad \forall \Delta t > 0.$$

and

Implicit negative derivative condition

$$u^{n+1} = u^n - \Delta t^2 \dot{\mathbf{G}}(u^{n+1}) \quad \implies \quad \|u^{n+1}\| \leq \|u^n\| \quad \forall \Delta t > 0.$$

for some convex functional $\|\cdot\|$.

A method given by

$$U = \mathbf{R}e u^n + \mathbf{P}U + \Delta t \mathbf{D}G(U) + \Delta t^2 \dot{\mathbf{D}}\dot{G}(U),$$

that satisfies the conditions

$$\mathbf{R}e \geq 0, \quad \mathbf{P} \geq 0, \quad \mathbf{D} \geq 0, \quad \dot{\mathbf{D}} \leq 0, \quad (\text{componentwise})$$

will preserve the strong stability property $\|u^{n+1}\| \leq \|u^n\|$ for any positive time-step $\Delta t > 0$.

These building blocks give diagonally implicit Runge–Kutta methods of order $p = 2, 3, 4$ that are unconditionally SSP.

S. Gottlieb, Z. Grant, J. Hu, R. Shu (2022).

Second order The one-stage, second order method is simply the implicit Taylor series method

$$u^{n+1} = u^n + \Delta t G(u^{n+1}) - \frac{1}{2} \Delta t^2 \dot{G}(u^{n+1}).$$

Third order

$$\mathbf{D} = \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}, \quad \dot{\mathbf{D}} = \begin{bmatrix} -\frac{1}{6} & 0 \\ 0 & -\frac{1}{3} \end{bmatrix},$$

$$\mathbf{P} = \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix}, \quad \mathbf{Re} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}.$$

Fourth order A five-stage, fourth order unconditionally SSP implicit two-derivative Runge–Kutta method is given by the Shu–Osher coefficients

$$\text{diag}(\mathbf{D}) = \begin{bmatrix} 0.660949255604937 \\ 0.242201390400848 \\ 1.137542996287740 \\ 0.191388711018110 \\ 0.625266691721946 \end{bmatrix}, \quad \text{diag}(\dot{\mathbf{D}}) = \begin{bmatrix} -0.177750705279127 \\ -0.354733903778084 \\ -0.403963513682271 \\ -0.161628266349058 \\ -0.218859021269943 \end{bmatrix},$$

$$P = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 \\ 0.084036809261019 & 0.915963190738981 & 0 & 0 & 0 \\ 0.001511648458457 & 0 & 0.090254853867587 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \end{bmatrix},$$

$$\mathbf{Re} = [1, 0, 0, 0, 0.908233497673956, 0]^T.$$

We don't often deal with problems $u_t = \mathbf{G}(u)$ that have these properties, but we do if we consider equations of the form

$$u_t = \mathbf{F}(u) + \mathbf{G}(u),$$

- \mathbf{F} satisfies a forward Euler condition where the time-step restriction is of a reasonable size (i.e. \mathbf{F} is non-stiff),
- \mathbf{G} satisfies an implicit Euler condition, with no time-step restriction.
- $\dot{\mathbf{G}}$ satisfies a negative derivative condition, with no time-step restriction.

Such equations are common in kinetic models; e.g. BGK, Broadwell, etc. These have a transport term \mathbf{F} and stiff collision term \mathbf{G} that satisfy these conditions.

An s -stage multi-derivatve IMEX method ($\mathcal{C} = r$)

$$u^{(i)} = r_i u^n + \sum_{j=1}^{i-1} p_{ij} u^{(j)} + \sum_{j=1}^{i-1} w_{ij} \left(u^{(j)} + \frac{\Delta t}{r} F(u^{(j)}) \right) \\ + \Delta t d_{ii} G(u^{(i)}) + \Delta t^2 \dot{d}_{ii} \dot{G}(u^{(i)}), \quad i = 1, \dots, s, \\ u^{n+1} = u^{(s)}.$$

In matrix form,

$$U = \mathbf{R}e u^n + \mathbf{P}U + \mathbf{W} \left(U + \frac{\Delta t}{r} F(U) \right) + \Delta t \mathbf{D}G(U) + \Delta t^2 \dot{\mathbf{D}}\dot{G}(U),$$

where \mathbf{P} , \mathbf{W} , and $\mathbf{R} = I - \mathbf{P} - \mathbf{W}$ are $s \times s$ matrices, \mathbf{D} and $\dot{\mathbf{D}}$ are $s \times s$ diagonal matrices, and e is a vector of ones.

SSP IMEX schemes with second derivative

Given operators F and G that satisfy for some $\|\cdot\|$:

Forward Euler condition on F

$$\|u + \Delta t F(u)\| \leq \|u\| \quad \text{for all } \Delta t \leq \Delta t_{\text{FE}},$$

for some $\Delta t_{\text{FE}} > 0$

Implicit Euler condition on G

$$u^{n+1} = u^n + \Delta t G(u^{n+1}) \quad \implies \quad \|u^{n+1}\| \leq \|u^n\| \quad \forall \Delta t > 0.$$

Implicit negative derivative condition on $\dot{G} = G'(u)G(u)$

$$u^{n+1} = u^n - \Delta t^2 \dot{G}(u^{n+1}) \quad \implies \quad \|u^{n+1}\| \leq \|u^n\| \quad \forall \Delta t > 0.$$

Under the conditions above on \mathbf{F} , \mathbf{G} , and $\dot{\mathbf{G}}$, if the method

$$U = \mathbf{R}e u^n + \mathbf{P}U + \mathbf{W} \left(U + \frac{\Delta t}{r} \mathbf{F}(U) \right) + \Delta t \mathbf{D} \mathbf{G}(U) + \Delta t^2 \dot{\mathbf{D}} \dot{\mathbf{G}}(U),$$

satisfies

$$\mathbf{R}e \geq 0, \quad \mathbf{P} \geq 0, \quad \mathbf{W} \geq 0, \quad \mathbf{D} \geq 0, \quad \dot{\mathbf{D}} \leq 0,$$

then it preserves the strong stability property $\|u^{n+1}\| \leq \|u^n\|$ under the time-step condition $\Delta t \leq r \Delta t_{\text{FE}}$.

We found such methods of up to order $p = 3$.

There is no time step restriction coming from \mathbf{G} , which is unusual for IMEX methods!

Second order method SSP IMEX schemes

We begin with a method that has Shu-Osher coefficients

$$\mathbf{W} = \begin{bmatrix} 0 & 0 & 0 \\ 1 & 0 & 0 \\ 0 & 1/2 & 0 \end{bmatrix}, \quad \mathbf{P} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 1/2 & 0 & 0 \end{bmatrix}, \quad \mathbf{Re} = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix},$$

and

$$\text{diag}(\mathbf{D}) = \begin{bmatrix} \frac{1}{2} \\ 0 \\ \frac{1}{2} \end{bmatrix}, \quad \text{diag}(\dot{\mathbf{D}}) = - \begin{bmatrix} 0 \\ \frac{1}{2} \\ 0 \end{bmatrix},$$

with $r = 1$.

Looser conditions

If we allow the final stage to be explicit and contain only \mathbf{F} terms, we get a larger SSP coefficient $\mathcal{C} = (1 + \sqrt{2})/2 > 1$:

$$y^{(1)} = u^n + \frac{1}{2 + \sqrt{2}} \Delta t \mathbf{G}(y^{(1)}) - \frac{1}{2 + \sqrt{2}} \Delta t^2 \dot{\mathbf{G}}(y^{(1)})$$

$$y^{(2)} = \left(y^{(1)} + \frac{\Delta t}{r} \mathbf{F}(y^{(1)}) \right)$$

$$y^{(3)} = \frac{6 - \sqrt{2}}{8} y^{(1)} + \frac{2 + \sqrt{2}}{8} \left(y^{(2)} + \frac{\Delta t}{r} \mathbf{F}(y^{(2)}) \right) + \frac{1}{\sqrt{2}} \Delta t \mathbf{G}(y^{(3)})$$

$$u^{n+1} = \left(1 - \frac{\sqrt{2}}{4} \right) y^{(3)} + \frac{2 + \sqrt{2}}{4(1 + \sqrt{2})} \left(y^{(3)} + \frac{\Delta t}{r} \mathbf{F}(y^{(3)}) \right).$$

Third order SSP IMEX schemes

This method has $r = 0.904402174130635$ with:

$$\text{diag}(\mathbf{D}) = \begin{pmatrix} 0 \\ 2 \\ 0.388820513661584 \\ 0.083529464436389 \\ 1.793313488277995 \\ 0 \end{pmatrix},$$
$$\text{diag}(\dot{\mathbf{D}}) = - \begin{pmatrix} 0.871358934880525 \\ 0.856842702601821 \\ 0 \\ 0 \\ 2 \\ 0.205134529930013 \end{pmatrix}.$$

Note that $d_{ii} + |\dot{d}_{ii}| > 0$ for each stage i .

Third order SSP IMEX scheme

$$\mathbf{W} = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0.058453072749259 & 0 & 0 & 0 & 0 & 0 \\ 0.764266518291495 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0.292520982667463 & 0 & 0 & 0 \\ 0.173788618990251 & 0 & 0 & 0.281050180194829 & 0 & 0 \\ 0.016811671845949 & 0 & 0 & 0.448630511341543 & 0 & 0 \end{pmatrix},$$

$$\mathbf{P} = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0.253395246357353 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0.235733481708505 & 0 & 0 & 0 & 0 \\ 0 & 0.123961833526104 & 0 & 0 & 0 & 0 \\ 0.409037644509411 & 0.136123556305509 & 0 & 0 & 0 & 0 \\ 0.203353399602184 & 0 & 0 & 0 & 0.331204417210324 & 0 \end{pmatrix},$$

$$\mathbf{Re} = \begin{pmatrix} 1 \\ 0.688151680893388 \\ 0 \\ 0.583517183806433 \\ 0 \\ 0 \end{pmatrix}.$$

Second order with $(k \geq 3)$ -step two-derivative two-stages

More steps allow larger SSP coefficients, e.g. a family of methods with $r = C = \frac{k-2}{k-1}$

$$\begin{aligned}y^{(1)} &= u^n - (k-1)\Delta t^2 \dot{\mathbf{G}}(y^{(1)}) \\y^{(2)} &= \frac{1}{k-1}u^{n-k+1} + \frac{k-2}{k-1} \left(y^{(1)} + \frac{\Delta t}{r} \mathbf{F}(y^{(1)}) \right) \\&\quad + k\Delta t \mathbf{G}(y^{(2)}) - k\Delta t^2 \dot{\mathbf{G}}(y^{(2)}) \\u^{n+1} &= \frac{1}{k-1}y^{(2)} + \frac{k-2}{k-1} \left(y^{(1)} + \frac{\Delta t}{r} \mathbf{F}(y^{(1)}) \right)\end{aligned}$$

Families of problems that satisfy these conditions

The negative derivative is naturally found as part of kinetic equations of the form

$$u_t = T(u) + \frac{1}{\varepsilon}Q(u)$$

where T satisfies a forward Euler condition

$$\|u^n + \Delta t T(u^n)\| \leq \|u^n\| \quad \forall \Delta t \leq \Delta t_F$$

and Q satisfies a forward Euler condition

$$\|u^n + \Delta t Q(u^n)\| \leq \|u^n\| \quad \forall \Delta t \leq \Delta t_Q$$

and a **negative derivative condition**

$$\|u^n - \Delta t^2 \dot{Q}(u^n)\| \leq \|u^n\| \quad \text{for } \Delta t \leq \Delta t_{\dot{Q}}$$

Numerical Results: Broadwell Model

The Broadwell model is a simple discrete velocity kinetic model:

$$\begin{cases} \partial_t f_+ + \partial_x f_+ = \frac{1}{\varepsilon}(f_0^2 - f_+ f_-), \\ \partial_t f_0 = -\frac{1}{\varepsilon}(f_0^2 - f_+ f_-), \\ \partial_t f_- - \partial_x f_- = \frac{1}{\varepsilon}(f_0^2 - f_+ f_-), \end{cases} \quad (1)$$

where $f_+ = f_+(t, x)$, $f_0 = f_0(t, x)$, and $f_- = f_-(t, x)$ denote the densities of particles with speed 1, 0, and -1 , respectively.

Define $f = (f_+, f_0, f_-)^T$,

$T(f) = (-\partial_x f_+, 0, \partial_x f_-)^T$,

$Q(f) = (f_0^2 - f_+ f_-, -(f_0^2 - f_+ f_-), f_0^2 - f_+ f_-)^T$

(these should be defined for the system after spatial discretization).



Numerical Results: Broadwell Model

Consider the Broadwell model (1) on the domain $x \in [0, 2]$ with periodic boundary condition and inconsistent initial data

$$\begin{aligned}f_+(0, \cdot) &= 1 + 0.2 \exp(0.3 \sin(\pi x)), & f_-(0, \cdot) &= \exp(0.2 \cos(2\pi x)), \\f_0(0, \cdot) &= \frac{1}{1 + 0.3 \sin(\pi x)}.\end{aligned}$$

We discretize in space by the fifth order finite volume WENO scheme, and the collision operator Q is evaluated pointwise on the Gauss quadrature points in each cell.

We fix the CFL number as $\Delta t = \frac{1}{2} \Delta x$, and solve (1) by the new second and third order methods up to final time $T = 0.1$.

Numerical Results: Broadwell Model

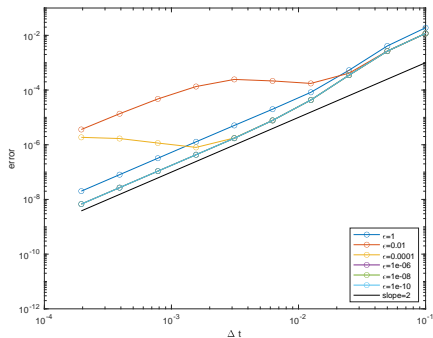


Figure: Accuracy test of the new second order IMEX scheme for the Broadwell model.

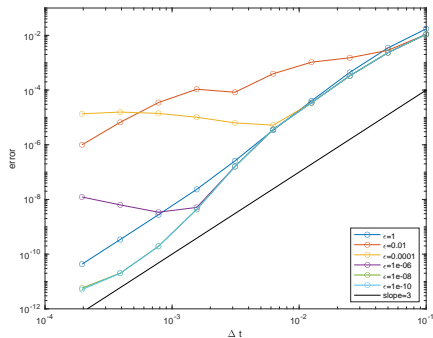


Figure: Accuracy test of the new third order IMEX scheme for the Broadwell model.

Example: The Bhatnagar-Gross-Krook (BGK) model

The negative derivative condition can be found in (part of) the BGK equation, a widely used kinetic model introduced to mimic the full Boltzmann equation:

$$\partial_t f + v \cdot \nabla_x f = \frac{1}{\varepsilon} (M - f), \quad x, v \in \mathbb{R}^d,$$

where $f = f(t, x, v)$ is the probability density function and M is the Maxwellian:

$$M(t, x, v) = \frac{\rho(t, x)}{(2\pi T(t, x))^{d/2}} \exp\left(-\frac{|v - u(t, x)|^2}{2T(t, x)}\right),$$

the density ρ , bulk velocity u and temperature T are given by the moments of f .

The Bhatnagar-Gross-Krook (BGK) model

We see that the BGK model has the property

$$\dot{Q}(f) = Q'Q = -Q(f).$$

So that if Q satisfies a forward Euler condition

$$\|u^n + \Delta t Q(u^n)\| \leq \|u^n\| \quad \forall \Delta t \leq \Delta t_Q$$

then \dot{Q} will automatically satisfy the **negative derivative condition**

$$\|u^n - \Delta t^2 \dot{Q}(u^n)\| = \|u^n + \Delta t^2 Q(u^n)\| \leq \|u^n\| \quad \text{for } \Delta t \leq \sqrt{\Delta t_Q}$$

Numerical Results: The BGK model

The BGK model is a widely used kinetic model introduced to mimic the full Boltzmann equation:

$$\partial_t f + v \cdot \nabla_x f = \frac{1}{\varepsilon} (M - f), \quad x, v \in \mathbb{R}^d, \quad (2)$$

where $f = f(t, x, v)$ is the probability density function and M is the so-called Maxwellian given by

$$M(t, x, v) = \frac{\rho(t, x)}{(2\pi T(t, x))^{d/2}} \exp\left(-\frac{|v - u(t, x)|^2}{2T(t, x)}\right), \quad (3)$$

where the density ρ , bulk velocity u and temperature T are given by the moments of f :

$$\rho = \int_{\mathbb{R}^d} f \, dv, \quad \rho u = \int_{\mathbb{R}^d} f v \, dv, \quad \frac{1}{2} \rho dT = \frac{1}{2} \int_{\mathbb{R}^d} f |v - u|^2 \, dv. \quad (4)$$

Numerical Results: The BGK model

Consider the 1D BGK model on the physical domain $x \in [0, 2]$ with periodic boundary condition, and inconsistent initial data given by

$$f(0, x, v) = 0.7M[\tilde{\rho}(x), \tilde{u}(x), \tilde{T}(x)](v) + 0.3M[\tilde{\rho}(x), -0.5\tilde{u}(x), \tilde{T}(x)](v)$$

with

$$\tilde{\rho}(x) = 1 + 0.2 \sin(2\pi x), \quad \tilde{u}(x) = 1, \quad \tilde{T}(x) = \frac{1}{1 + 0.2 \sin(\pi x)}.$$

The velocity domain is truncated into $[-v_{max}, v_{max}]$ with $v_{max} = 15$ and discretized with $N_v = 150$ grid points, and the physical space is discretized in the same way as for the Broadwell model.

We fix the the CFL number as $\Delta t = \frac{1}{2} \frac{\Delta x}{v_{max}}$.

Numerical Results: The BGK model

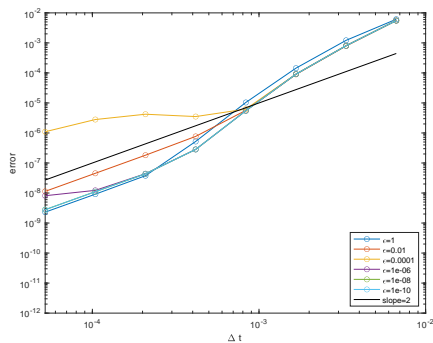


Figure: Accuracy test of the new second order IMEX scheme for the BGK model.

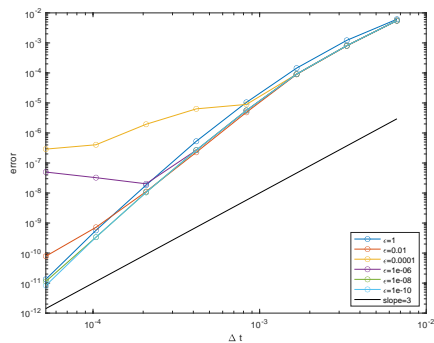


Figure: Accuracy test of the new third order IMEX scheme for the BGK model.

Additive SSP methods

- Additive methods handle T and Q with different coefficients, and ensure the SSP property will be preserved for the combination.
- We want methods that will handle T with downwinding because it is a transport term, so that T satisfies a forward Euler condition, and has a downwind version \tilde{T} .
- Because Q satisfies a forward Euler condition and a negative derivative condition, we want the additive method to handle Q with an implicit term only and an implicit negative derivative term only.

Additive SSP methods

If T and Q satisfy the conditions above, then an additive method of the form

$$\begin{aligned} u^{(i)} &= r_i u^n + \sum_{j=1}^{i-1} p_{ij} u^{(j)} + \sum_{j=1}^{i-1} w_{ij} \left(u^{(j)} + \frac{\Delta t}{r} T(u^{(j)}) \right) \\ &\quad + \sum_{j=1}^{i-1} \tilde{w}_{ij} \left(u^{(j)} - \frac{\Delta t}{r} T(u^{(j)}) \right) \\ &\quad + \frac{1}{\varepsilon} \Delta t d_{ii} Q(u^{(i)}) + \frac{1}{\varepsilon^2} \Delta t^2 \dot{d}_{ii} \dot{Q}(u^{(i)}), \quad i = 1, \dots, s, \\ u^{n+1} &= u^{(s)} \end{aligned}$$

is SSP with $\mathcal{C} = r$ if $r_i, p_{ij}, w_{ij}, \tilde{w}_{ij}$ and d_{ii} are all non-negative and \dot{d}_{ii} is non-positive.

In matrix form

$$U = \mathbf{R}e u^n + \mathbf{P}U + \mathbf{W} \left(U + \frac{\Delta t}{r} T(U) \right) + \tilde{\mathbf{W}} \left(U - \frac{\Delta t}{r} \tilde{T}(U) \right) \\ + \frac{\Delta t}{\epsilon} \mathbf{D}Q(U) + \frac{\Delta t^2}{\epsilon} \dot{\mathbf{D}}\dot{Q}(U),$$


where \mathbf{P} , \mathbf{W} , $\tilde{\mathbf{W}}$, and $\mathbf{R} = \mathbf{I} - \mathbf{P} - \mathbf{W} - \tilde{\mathbf{W}}$ are $s \times s$ matrices,
 r_i is the i -th row sum of \mathbf{R} ,
 \mathbf{D} and $\dot{\mathbf{D}}$ are $s \times s$ diagonal matrices,
and e is a vector of ones.

If all elements of $\mathbf{R}e$, \mathbf{P} , \mathbf{W} , $\tilde{\mathbf{W}}$ and \mathbf{D} are non-negative, and the elements of $\dot{\mathbf{D}}$ are non-positive, then this method is SSP with $\mathcal{C} = r$.

Can we find additive SSP methods with arbitrarily large r ?

Ongoing work by Andrew Christlieb, **Sining Gong**, Sigal Gottlieb, and Zack Grant found an unconditional additive method using downwinding for T and the second derivative for Q :

$$\begin{aligned}u^{(i)} &= r_i u^n + \sum_{j=1}^s p_{ij} u^{(j)} + \sum_{j=1}^s w_{ij} \left(u^{(j)} + \frac{\Delta t}{r} T(u^{(j)}) \right) \\ &\quad + \sum_{j=1}^s \tilde{w}_{ij} \left(u^{(j)} - \frac{\Delta t}{r} \tilde{T}(u^{(j)}) \right) \\ &\quad + \frac{1}{\varepsilon} \Delta t d_{ii} Q(u^{(i)}) + \frac{1}{\varepsilon^2} \Delta t^2 \dot{d}_{ii} \dot{Q}(u^{(i)}), \quad i = 1, \dots, s\end{aligned}$$

$$u^{n+1} = u^{(s)}.$$


Unconditional additive second order SSP methods

The coefficients are given by:

$$\mathbf{W} = \begin{bmatrix} 0 & 0 & \frac{2}{4+r} \\ \frac{1}{2r^2} & \frac{5}{13} & \frac{3}{8r} - \frac{3}{4r^2} \\ 0 & \frac{8}{13} & 0 \end{bmatrix}, \quad \tilde{\mathbf{W}} = \begin{bmatrix} \frac{2+r}{4+r} & 0 & 0 \\ 0 & 0 & \frac{8}{13} - \frac{1}{4r^2} - \frac{3}{8r} \\ 0 & 0 & \frac{5}{13} \end{bmatrix}$$

$$\text{diag}(\mathbf{D}) = \begin{bmatrix} \frac{2}{4+r} \\ 0 \\ 0 \end{bmatrix}, \quad \text{diag}(\dot{\mathbf{D}}) = \begin{bmatrix} 0 \\ -\frac{3}{8r^2} \\ -\frac{3}{8r^2} \end{bmatrix}, \quad \mathbf{Re} = \begin{bmatrix} 0 \\ \frac{1}{2r^2} \\ 0 \end{bmatrix}, \quad \mathbf{P} = \mathbf{0},$$

Note that when $r \geq 2$, all the coefficients here are positive.

Unconditional Additive SSP methods

Given

$$u_t = T(u) + \frac{1}{\epsilon} Q(u),$$

where

- T satisfies a forward Euler condition for some $\Delta t_{FE} > 0$.
- \tilde{T} satisfies a backward Euler condition for some $\Delta t_{DW} > 0$.
- $\frac{1}{\epsilon} Q$ satisfies an implicit Euler condition, with no time-step restriction.
- $\frac{1}{\epsilon^2} \dot{Q}$ satisfies a negative derivative condition, with no time-step restriction.

We found a family of second order SSP methods with arbitrarily large \mathcal{C} !

Downwinding for the transport term and negative derivative conditions for the collision term give good* additive SSP methods!

Conclusions

SSP properties depend on properties of the underlying problem and discretization, so that it is helpful to have different methods

- Optimal SSP Runge–Kutta methods with forward Euler and second derivative building blocks are in [A.J. Christlieb, S. Gottlieb, Z. Grant, and D. C. Seal, JSC \(2016\)](#)
- Optimal SSP Runge–Kutta methods with forward Euler and Taylor series building blocks are in [Z. Grant, S. Gottlieb, D.C. Seal, CAMC \(2019\)](#)
- Optimal SSP Runge–Kutta methods with forward Euler and negative derivative building blocks, [S. Gottlieb, Z. Grant, J. Hu, R. Shu, SINUM \(2022\)](#)
- A review as well as an extension to multi-step, multi-stage, two-derivative methods in *Applied Mathematical Modelling* Volume 149, January 2026.
- **Ongoing** Unconditional SSP Runge–Kutta methods using downwinding for the transport term and the implicit negative derivative building blocks for the collision term [A. Christlieb, S. Gong, S. Gottlieb, Z. Grant](#)

Thank You!