

Well-posedness and invariance of the Gibbs measures for periodic generalized Korteweg-de Vries equations

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The generalized KdV equations

Periodic generalized Korteweg-de Vries equations:

$$(gKdV) \quad \begin{cases} \partial_t u + \partial_x^3 u = \pm \partial_x(u^k), \\ u|_{t=0} = u_0, \end{cases} \quad (t, x) \in \mathbb{R} \times \mathbb{T}$$

■ Dispersive equations

■ Hamiltonian equations: $(gKdV) \iff \partial_t u = \partial_x \frac{\delta H}{\delta u}$

$$H(u) = \frac{1}{2} \int_{\mathbb{T}} (\partial_x u)^2 dx \pm \frac{1}{k+1} \int_{\mathbb{T}} u^{k+1} dx$$

■ conserved quantities: **mean** $\int_{\mathbb{T}} u dx$, **mass** $\int_{\mathbb{T}} u^2 dx$, **Hamiltonian** $H(u)$

AIM: Construct global-in-time dynamics in the support of the **Gibbs measure** μ

$$d\mu = Z^{-1} e^{-H(u)} du$$

and show invariance of μ under the flow

Hamiltonian systems

Finite dimensional Hamiltonian system: $(p, q) \in \mathbb{R}^{2n}$ solve

$$\dot{p} = \frac{\partial H}{\partial q}, \quad \dot{q} = -\frac{\partial H}{\partial p}$$

with Hamiltonian $H = H(p, q)$.

▷ **Liouville's theorem + conservation of Hamiltonian** \implies Gibbs measure μ

$$d\mu = Z^{-1} \exp(-H(p, q)) dp dq$$

is **invariant** under the dynamics.

▷ Similarly for **gKdV**, the Gibbs measure formally defined as

$$d\mu \text{ " = " } Z^{-1} e^{-H(u)} du$$

is expected to be **invariant**:

$$\mu(\Psi(-t)A) = \mu(A), \quad t \in \mathbb{R}, \quad A \text{ measurable}$$

Why do we care about Gibbs measures?

1 Invariance of Gibbs measure \implies **typical** behaviour of solutions

- Poincaré recurrence theorem
- Furstenberg multiple recurrence theorem

2 **Local** well-posedness (LWP) + invariance \implies **global** well-posedness

- Bourgain's invariant measure argument ('94)
- Main idea: use invariance of Gibbs measure as a *substitute* for a conservation law

Problem 1: How do we make sense of the measure μ

$$d\mu \stackrel{?}{=} Z^{-1} e^{-H(u)} du \quad ?$$

Problem 2: Need local well-posedness at low regularity (in the support of μ)!

Gaussian measure

Want to make sense of the Gibbs measure μ as a **weighted Gaussian measure**

$$d\mu = Z^{-1} e^{\mp \frac{1}{k+1} \int_{\mathbb{T}} u^{k+1} dx} e^{-\frac{1}{2} \int_{\mathbb{T}} (\partial_x u)^2 dx} du$$

The **Gaussian measure** ρ can be defined through the density

$$d\rho = Z_0^{-1} e^{-\frac{1}{2} \int_{\mathbb{T}} (\partial_x u)^2 dx} du \quad (= \lim_{N \rightarrow \infty} d\rho_N)$$

or equivalently, the induced probability measure under the map

$$\omega \in \Omega \mapsto u^\omega(x) = \sum_{n \neq 0} \frac{g_n(\omega)}{|n|} e^{inx}$$

where $\{g_n\}_{n \neq 0}$ are independent \mathbb{C} -valued standard Gaussian r.v.'s with $g_{-n} = \bar{g}_n$

▷ ρ is a **probability measure** on $H^s(\mathbb{T}) \iff s < \frac{1}{2}$

$$\mathbb{E}[\|u^\omega\|_{H^s}^2] = \mathbb{E}\left[\sum_{n \neq 0} \frac{|n|^{2s} |g_n(\omega)|^2}{|n|^2}\right] = \sum_{n \neq 0} |n|^{2(s-1)} \underbrace{\mathbb{E}[|g_n(\omega)|^2]}_{=1}$$

Gibbs measure

- **defocusing case:** '+' in gKdV and k odd

$$d\mu = Z^{-1} e^{-\frac{1}{k+1} \int_{\mathbb{T}} u^{k+1} dx} d\rho$$

▷ $u \in \text{supp } \rho \implies u \in L^{k+1}(\mathbb{T})$ (Sobolev embedding)

- **non-defocusing case:** '-' in gKdV or k even

Lebowitz-Rose-Speer '88:
$$d\mu = Z^{-1} \mathbb{1}_{\{\|u\|_{L^2} \leq R\}} e^{\mp \frac{1}{k+1} \int_{\mathbb{T}} u^{k+1} dx} d\rho$$

Theorem (Lebowitz-Rose-Speer '88, Bourgain '94, Oh-Sosoe-Tolomeo '21)

Let $k \geq 2$, $R > 0$ and $Z_{k,R} = \left\| \mathbb{1}_{\{\|u\|_{L^2} \leq R\}} e^{\mp \frac{1}{k+1} \int_{\mathbb{T}} u^{k+1} dx} \right\|_{L^1(\rho)}$.

- $2 \leq k \leq 4$: $Z_{k,R} < \infty$ for any $R > 0$
- $k = 5$: $Z_{5,R} < \infty$ if $R \leq \|Q\|_{L^2(\mathbb{R})}$; $Z_{5,R} = \infty$ if $R > \|Q\|_{L^2(\mathbb{R})}$

▷ **Also:** μ is well approximated by μ_N defined by

$$d\mu_N = Z^{-1} \mathbb{1}_{\{\|P_{\leq N} u\|_{L^2} \leq R\}} e^{\mp \frac{1}{k+1} \int_{\mathbb{T}} (P_{\leq N} u)^{k+1} dx} d\rho$$

Bourgain '94: Construct **local-in-time** solutions and exploit the (formal) invariance of the Gibbs measure μ as a substitute for a conservation law

▷ **General strategy:**

show **local well-posedness** in the support of μ
+
Bourgain's invariant measure argument

▷ **Main difficulty:** The support of the Gibbs measure μ is rough

$$\text{supp } \mu \subset H^{\frac{1}{2}-\varepsilon}(\mathbb{T}) \setminus H^{\frac{1}{2}}(\mathbb{T}), \quad \varepsilon > 0$$

Known results

KdV ($k = 2$)	Bourgain '93	GWP in $L^2(\mathbb{T})$
mKdV ($k = 3$)	Bourgain '94	a.s. GWP in $H^{\frac{1}{2}-}(\mathbb{T}) \cap \mathcal{FL}^{1-, \infty}(\mathbb{T})$
	Bourgain '93, Staffilani '97	
	Colliander-Keel-Staffilani -Takaoka-Tao '04	LWP in $H^{\frac{1}{2}}(\mathbb{T})$ GWP in $H^{\frac{13}{14} - \frac{2}{7(k-1)}}(\mathbb{T})$
gKdV ($k \geq 4$)		C^k failure in $H^{\frac{1}{2}-}(\mathbb{T})$
	Richards '16 ($k = 4$)	a.s. LWP in $H^{\frac{1}{2}-}(\mathbb{T})$, a.s. GWP invariance under \mathcal{G} -gKdV
	Oh-Richards-Thomann '16	a.s. GWP in $H^{\frac{1}{2}-}(\mathbb{T})$ (no uniqueness) "mild" invariance [†] under gKdV

† following a compactness argument due to Burq-Thomann-Tzvetkov '18, they show that $\mathcal{L}(u(t)) = \mu$ for all $t \in \mathbb{R}$

Problem: $\text{supp } \mu \subset H^{\frac{1}{2}-}(\mathbb{T})$ where gKdV is mildly ill-posed

Alternative 1: Show *probabilistic* well-posedness in $H^{\frac{1}{2}-}(\mathbb{T})$

- Consider random initial data $u^\omega(x) = \sum_{n \neq 0} \frac{g_n(\omega)}{|n|} e^{inx}$

Alternative 2: Show *deterministic* local well-posedness in different spaces

- **Fourier-Lebesgue spaces** $\mathcal{FL}^{s,p}(\mathbb{T})$:

$$\|f\|_{\mathcal{FL}^{s,p}} = \|\langle n \rangle^s \widehat{f}(n)\|_{\ell_n^p}, \quad \langle n \rangle = (1 + n^2)^{\frac{1}{2}}$$

- $\mathcal{FL}^{s,2}(\mathbb{T}) = H^s(\mathbb{T})$

- $\text{supp } \mu \subset \mathcal{FL}^{s,p}(\mathbb{T}) \iff s < 1 - \frac{1}{p}$

$$\mathbb{E}[\|u^\omega\|_{\mathcal{FL}^{s,p}}^p] \sim \sum_{n \neq 0} |n|^{(s-1)p} < \infty \iff (s-1)p < -1$$

- $\mathcal{FL}^{s,p}(\mathbb{T}) \hookrightarrow L^{k+1}(\mathbb{T}), s > 1 - \frac{1}{p} - \frac{1}{k+1}$

▷ **Gauged gKdV equation:**

$$(\mathcal{G}\text{-gKdV}) \quad \begin{cases} \partial_t u + \partial_x^3 u = \pm \partial_x (u^k - k \mathbf{P}_0(u^{k-1})u) \\ u|_{t=0} = u_0 \end{cases}$$

▷ **Gauge transform:** $\mathcal{G}[u](t, x) = u\left(t, x \mp k \int_0^t \mathbf{P}_0(u^{k-1}(t')) dt'\right)$

Theorem (Local well-posedness of \mathcal{G} -gKdV) (C.-Kishimoto '21)

Let $k \geq 3$ and $2 < p < \infty$. Then, there exists $\varepsilon_0 = \varepsilon_0(k, p) \in (0, \frac{1}{k+1})$ such that \mathcal{G} -gKdV is locally well-posed in $\mathcal{FL}^{s,p}(\mathbb{T})$ for any $s > \max(\frac{1}{2}, 1 - \frac{1}{p} - \varepsilon_0)$.

- Fully deterministic result
- **Key ingredients:** multilinear Strichartz estimates adapted to the Fourier-Lebesgue setting

Main results

Theorem (a.s. GWP of \mathcal{G} -gKdV and invariance of μ) (C.-Kishimoto '21)

Let $k \geq 3$ and assume that the measure μ is well-defined. Then, there exists $\Sigma \subset \bigcap_{s < 1 - \frac{1}{p}} \mathcal{F}L^{s,p}(\mathbb{T})$ of full μ -measure s.t.

- 1 For every $u_0 \in \Sigma$, there exists a **unique global-in-time** solution $u \in \bigcap_{s < 1 - \frac{1}{p}} C(\mathbb{R}; \mathcal{F}L^{s,p}(\mathbb{T}))$ of \mathcal{G} -gKdV
- 2 The solution map $\Phi(t)$ for \mathcal{G} -gKdV is μ -measurable, $\Phi(t)\Sigma = \Sigma$, and $\Phi(t+s) = \Phi(t)\Phi(s)$
- 3 The Gibbs measure μ is **invariant** under the flow of \mathcal{G} -gKdV

By inverting the gauge transform and using the invariance of the Gibbs measure under translations, we obtain the following:

Theorem (C.-Kishimoto '21)

Under the assumptions of the previous theorem, the (ungauged) gKdV equation is a.s. globally well-posed and the Gibbs measure is invariant under its flow.

Linear propagator: $\mathcal{F}_x(S(t)u_0)(n) = e^{itn^3}\hat{u}_0(n)$

Duhamel formulation:

$$u(t) = S(t)u_0 \pm \int_0^t S(t-t') \underbrace{\partial_x(u^k - k\mathbf{P}_0(u^{k-1})u)}_{\mathcal{N}[u](t')} dt'$$

Interaction representation: $v(t) = S(-t)u(t)$ solves

$$\hat{v}(t, n) = \hat{u}_0(n) \pm \int_0^t \sum_{n=n_1+\dots+n_k} e^{-it'\phi_k(\bar{n})} i n \prod_{j=1}^k \hat{v}(t', n_j) dt'$$

Phase function: $\phi_k(\bar{n}) = n^3 - n_1^3 - \dots - n_k^3$

▷ **Lack of smoothing** effect in the periodic setting forces us to use the phase function to control the **derivative n**

GOOD: ← non-resonant

BAD: ← resonant

Fourier restriction norm method

Duhamel formulation:
$$u(t) = \underbrace{S(t)u_0}_{\text{linear solution}} \pm \int_0^t S(t-t')\mathcal{N}[u](t') dt'$$

Fourier restriction spaces $X_{p,q}^{s,b}$: (Bourgain '93, Grünrock '04, G.-Herr '08)

$$\|u\|_{X_{p,q}^{s,b}} := \left\| \langle n \rangle^s \left\| \langle \tau - n^3 \rangle^b \widehat{u}(\tau, n) \right\|_{L_\tau^q} \right\|_{\ell_n^p}$$

$\triangleright X_{p,2}^{s,\frac{1}{2}} \cap X_{p,1}^{s,0} \quad (\hookrightarrow C(\mathbb{R}; FL^{s,p}(\mathbb{T})))$

Main nonlinear estimate:
$$\|\mathcal{N}[u_1, \dots, u_k]\|_{X_{p,2}^{s,-\frac{1}{2}} \cap X_{p,1}^{s,-1}} \lesssim \prod_{j=1}^k \|u_j\|_{X_{p,2}^{s,\frac{1}{2}} \cap X_{p,1}^{s,0}}$$

Main ideas: Use smoothing in time to help control the spatial derivative

$$\underbrace{|n^3 - n_1^3 - \dots - n_k^3|}_{\phi_k} = |(\tau - n^3) - (\tau_1 - n_1^3) - \dots - (\tau_k - n_k^3)| \lesssim \max(|\tau - n^3|, |\tau_j - n_j^3|)$$

and exploit *dispersion* through Strichartz estimates

Phase function ϕ_k :

$$|\phi_k| \lesssim \max(|\tau - n^3|, |\tau_j - n_j^3|)$$

KdV $\phi_2 = n^3 - n_1^3 - n_2^3 \quad n = \underline{n_1 + n_2} \quad 3nn_1n_2$

mKdV $\phi_3 = n^3 - n_1^3 - n_2^3 - n_3^3 \quad n = \underline{n_1 + n_2 + n_3} \quad 3(n - n_1)(n - n_2)(n - n_3)$

gKdV ($k \geq 4$) $\phi_k = n^3 - n_1^3 - \dots - n_k^3 \quad n = \underline{n_1 + \dots + n_k} \quad ??$

Problem: no known factorization for ϕ_k when $k \geq 4$! (assume $|n_1| \geq \dots \geq |n_k|$)

Idea: compare ϕ_k with ϕ_2 and ϕ_3

$$\begin{aligned}\phi_k &= n^3 - n_1^3 - (n - n_1)^3 + (n - n_1)^3 - n_2^3 - \dots - n_k^3 \\ &= 3nn_1(n - n_1) + \mathcal{O}(|n_2n_3| \max(|n_4|, |n - n_1|))\end{aligned}$$

$$\begin{aligned}\phi_k &= n^3 - n_1^3 - n_2^3 - (n - n_1 - n_2)^3 + (n_3 + \dots + n_k)^3 - n_3^3 - \dots - n_k^3 \\ &= 3(n_1 + n_2)(n - n_1)(n - n_2) + \mathcal{O}(|n_3n_4| \max(|n_5|, |n - n_1 - n_2|))\end{aligned}$$

But first need to consider resonances: $\phi_2 = nn_1(n - n_1) = 0$
 $\phi_3 = (n_1 + n_2)(n - n_1)(n - n_2) = 0$

Gauged nonlinearity: $\mathcal{N}[u_1, \dots, u_k] = \partial_x \left(u_1 \cdots u_k - \sum_{j=1}^k \mathbf{P}_0 \left(\prod_{i \neq j} u_i \right) u_j \right)$

$$\implies \mathcal{F}_x \mathcal{N}[u_1, \dots, u_k](n) = \sum_{\substack{n=n_1+\dots+n_k \\ n \cdot n_1 \cdots n_k \neq 0}} \left(1 - \sum_{j=1}^k \mathbb{1}_{\{n=n_j\}} \right) i n \hat{u}_1(n_1) \cdots \hat{u}_k(n_k)$$

$$\mathcal{F}_x \mathcal{N}_0[u_1, \dots, u_k](n) = \sum_{\substack{n=n_1+\dots+n_k \\ n \cdot n_1 \cdots n_k \neq 0}} \mathbb{1}_{\bigcap_{j=1}^k \{n \neq n_j\}} i n \hat{u}_1(n_1) \cdots \hat{u}_k(n_k)$$

$$\mathcal{F}_x \mathcal{R}[u_1, \dots, u_k](n) = \sum_{\substack{n=n_1+\dots+n_k \\ n \cdot n_1 \cdots n_k \neq 0}} \sum_{\substack{J \subset \{1, \dots, k\} \\ |J| \geq 2}} (-1)^{|J|} \mathbb{1}_{\bigcap_{j \in J} \{n=n_j\}} i n \hat{u}_1(n_1) \cdots \hat{u}_k(n_k)$$

Then, $\boxed{\mathcal{N} = \mathcal{N}_0 + \mathcal{R}}$

Estimate for resonant part \mathcal{R} : Consider, e.g., the case $n = n_1 = n_2$

$$\langle n \rangle^s n \hat{u}_1(n) \hat{u}_2(n) \sum_{-n=n_3+\dots+n_k} \hat{u}_3(n_3) \cdots \hat{u}_k(n_k)$$

▷ Assuming $|n_3| \geq \dots \geq |n_k|$, we can use $\hat{u}_1, \hat{u}_2, \hat{u}_3$ to control the derivative

Estimate for the non-resonant part \mathcal{N}_0 :

$$\begin{aligned}\phi_k &= 3nn_1(n - n_1) + \mathcal{O}(|n_2n_3| \max(|n_4|, |n - n_1|)) \\ &= 3(n_1 + n_2)(n - n_1)(n - n_2) + \mathcal{O}(|n_3n_4| \max(|n_5|, |n - n_1 - n_2|))\end{aligned}$$

Further splitting:

$$\begin{aligned}\text{if } n_1 + n_2 &\neq 0 && \rightarrow \mathcal{N}_1 \\ \text{if } n_1 + n_2 = 0, n_3 + n_4 &\neq 0 && \rightarrow \mathcal{N}_3 \\ \text{if } n_1 + n_2 = n_3 + n_4 = 0, n_5 + n_6 &\neq 0 && \rightarrow \mathcal{N}_5 \\ \vdots &&& \vdots\end{aligned}$$

Then, $\mathcal{N}_0 = \mathcal{N}_1 + \mathcal{N}_3 + \dots$

For $1 \leq \alpha \leq k$ odd, then $n = n_\alpha + \dots + n_k$ with $n_\alpha + n_{\alpha+1} \neq 0$, so we consider

$$\begin{aligned}\phi_k &= 3nn_\alpha(n - n_\alpha) + \dots \\ &= 3(n_\alpha + n_{\alpha+1})(n - n_\alpha)(n - n_{\alpha+1}) + \dots\end{aligned}$$

Example: $|n| \sim |n_1| \gg |n_2| \geq \dots \geq |n_k|$

$$\phi_k = 3nn_1(n - n_1) + \mathcal{O}(|n_2|, |n_3|, \max(|n_4|, |n - n_1|))$$

• **Case 1:** $|\phi_k| \gtrsim |nn_1(n - n_1)|$ ($\lesssim \max\{\langle \tau - n^3 \rangle, \langle \tau_j - n_j^3 \rangle\} =: \sigma_{\max}$)

▷ Use 'time derivatives' to control spatial derivatives!

• **Case 2:** $|\phi_k| \ll |nn_1(n - n_1)|$

$$\implies |n_1|^2 |n - n_1| \lesssim |n_2 n_3| \max(|n_4|, |n - n_1|) \lesssim |n_2 n_3 n_4|$$

$$\implies \langle n \rangle^s |n| \lesssim \frac{\langle n_1 \rangle^s |n_2 n_3 n_4|^{\frac{1}{2}}}{|n - n_1|^{\frac{1}{2}}}$$

▷ **Key ingredients:** bilinear ($\sim L^4$) and trilinear ($\sim L^6$) Strichartz estimates adapted to the Fourier-Lebesgue setting!

(Linear) Strichartz estimates (Bourgain '93):

$$\|u\|_{L^4_{t,x}} \lesssim \|u\|_{X_{2,2}^{0,\frac{1}{3}}}, \quad \|u\|_{L^6_{t,x}} \lesssim \|u\|_{X_{2,2}^{0+, \frac{1}{2}}}$$

Multilinear Strichartz estimates (C.-Kishimoto '21)

For any $2 \leq p < \infty$, $b > \max(\frac{1}{3}, \frac{3p-2}{8p})$, we have

$$\|\mathbf{P}_{\neq 0}(\mathbf{P}_{\neq 0}u_1 \mathbf{P}_{\neq 0}u_2)\|_{X_{p,2}^{0,0}} \lesssim \|u_1\|_{X_{p,2}^{0,b}} \|u_2\|_{X_{2,2}^{0,b}} \quad (p = 2: \sim L^4)$$

$$\|\mathbf{P}_{\phi_3 \neq 0}(u_1 u_2 u_3)\|_{X_{p,2}^{0,0}} \lesssim \|u_1\|_{X_{p,2}^{0+, \frac{1}{2}}} \|u_2\|_{X_{2,2}^{0+, \frac{1}{2}}} \|u_3\|_{X_{2,2}^{0+, \frac{1}{2}}} \quad (p = 2: \sim L^6)$$

Bilinear estimate for $p \geq 2$:

Bilinear estimate (C.-Kishimoto '21)

For $2 \leq p < \infty$ and $b > \max(\frac{1}{3}, \frac{3p-2}{8p})$, we have

$$\|\mathbf{P}_{\neq 0}(\mathbf{P}_{\neq 0}u_1 \mathbf{P}_{\neq 0}u_2)\|_{X_{p,2}^{0, -\frac{1}{2}+}} \lesssim \|u_1\|_{X_{p,2}^{0,b}} \|u_2\|_{X_{p',2}^{0,0}}$$

$$|n| \sim |n_1| \gg |n_2| \geq \dots |n_k| \text{ and } |\phi_k| \ll |nn_1(n - n_1)|$$

$$\|\mathbf{P}_{\neq 0}(\mathbf{P}_{\neq 0} u_1 \mathbf{P}_{\neq 0} u_2)\|_{X_{p,2}^{0, -\frac{1}{2}+}} \lesssim \|u_1\|_{X_{p,2}^{0,b}} \|u_2\|_{X_{p',2}^{0,0}}$$

$$\langle n \rangle^s |n| \lesssim \frac{\langle n_1 \rangle^s |n_2 n_3 n_4|^{1/2}}{\langle n - n_1 \rangle^{1/2}}$$

$$\|\partial_x^s u_1 \cdot \partial_x^{-\frac{1}{2}} (\partial_x^{\frac{1}{2}} u_2 \cdot \partial_x^{\frac{1}{2}} u_3 \cdot \partial_x^{\frac{1}{2}} u_4 \prod_{j=5}^k u_j)\|_{X_{p,2}^{0, -\frac{1}{2}+}}$$

bilinear est.

$$\lesssim \|u_1\|_{X_{p,2}^{s, \frac{1}{2}}} \|\partial_x^{\frac{1}{2}} u_2 \cdot \partial_x^{\frac{1}{2}} u_3 \cdot \partial_x^{\frac{1}{2}} u_4 \prod_{j=5}^k u_j\|_{X_{p',2}^{-\frac{1}{2}, 0}}$$

$$\lesssim \|u_1\|_{X_{p,2}^{s, \frac{1}{2}}} \|\partial_x^{\frac{1}{2}} u_2 \cdot \partial_x^{\frac{1}{2}} u_3 \cdot \partial_x^{\frac{1}{2}} u_4\|_{X_{p,2}^{0,0}} \prod_{j=5}^k \|u_j\|_{X_{1,1}^{0,0}}$$

trilinear
Strichartz est.

$$\lesssim \|u_1\|_{X_{p,2}^{s, \frac{1}{2}}} \|u_2\|_{X_{p,2}^{\frac{1}{2}+, \frac{1}{2}}} \|u_3\|_{X_{2,2}^{\frac{1}{2}+, \frac{1}{2}}} \|u_4\|_{X_{2,2}^{\frac{1}{2}+, \frac{1}{2}}} \prod_{j=5}^k \|u_j\|_{X_{1,1}^{0,0}}$$

□

Invariance of the Gibbs measure for \mathcal{G} -gKdV

Bourgain's invariant measure argument:

- $\Phi(t)u_0 = u(t), \forall t \in \mathbb{R}$
- $\mu(\Phi(-t)A) = \mu(A), \forall t \in \mathbb{R}, A$
meas.

$$\begin{array}{ccc} u_N & \longrightarrow & u \\ \text{Invariance} \parallel & & \vdots \\ \mu_N & \longrightarrow & \mu \end{array}$$

Study **truncated gauged gKdV equation**:

$$(\mathbf{P}(f) = f - k\mathbf{P}_0(f))$$

$$(\mathcal{G}\text{-gKdV}_N) \quad \begin{cases} \partial_t u_N + \partial_x^3 u_N = \mathbf{P}_{\leq N} \partial_x [\mathbf{P}((\mathbf{P}_{\leq N} u_N)^{k-1}) \mathbf{P}_{\leq N} u_N] \\ u_N|_{t=0} = u_0 \end{cases}$$

$$u_{\text{high}} = \mathbf{P}_{>N} u_N \rightarrow \text{linear solution}$$

$$u_{\text{low}} = \mathbf{P}_{\leq N} u_N \rightarrow \text{finite system of ODEs}$$

Gibbs measure $\mu_N = \tilde{\mu}_N \otimes \rho_N^\perp$ associated with \mathcal{G} -gKdV $_N$

$$d\tilde{\mu}_N = \tilde{Z}_N^{-1} \mathbb{1}_{\{\|\mathbf{P}_{\leq N} u_N\|_{L^2} \leq R\}} e^{\mp \frac{1}{k+1} \int_{\mathbb{T}} (\mathbf{P}_{\leq N} u_N)^{k+1} dx} \underbrace{e^{-\frac{1}{2} \int_{\mathbb{T}} (\mathbf{P}_{\leq N} u_N)^2 dx} d(\mathbf{P}_{\leq N} u_N)}_{d\rho_N(u_N)}$$

is invariant under the flow!

What about gKdV?

Recall the inverse “**gauge transform**”

$$\mathcal{G}^{-1}[u](t, x) = u\left(t, x \pm k \int_0^t \mathbf{P}_0(u^{k-1}(t')) dt'\right)$$

▷ a.s. **GWP of gKdV** follows from a.s. GWP of \mathcal{G} -gKdV:

$$u \text{ solves } \mathcal{G}\text{-gKdV} \implies \mathcal{G}^{-1}[u] \text{ solves gKdV}$$

▷ **Invariance** $\mu(\Psi(-t)A) = \mu(A)$

$$\Psi(t)u_0 = \Phi(t)u_0\left(\cdot \mp k \int_0^t \int_{\mathbb{T}} \Phi(t')u_0^{k-1} dt'\right) = \tau_{\beta(t, u_0)}\Phi(t)u_0$$

■ **Problem:** cannot use invariance of μ under spatial translations

■ **Solution:** average over all translations!

$$\mu(\Psi(-t)A) = \int_{\mathbb{T}} \mu(\tau_y \Psi(-t)A) dm(y) = \int_{\mathbb{T}} \int_{\Sigma} \mathbb{1}_A(\tau_{y+\beta(t, u_0)}\Phi(t)u_0) d\mu(u_0) dm(y)$$

Thank you for your attention!